

THE PROFILES OF THE 3–12 MICRON POLYCYCLIC AROMATIC HYDROCARBON FEATURES¹

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ABSTRACT

We present spectra of the 3.3 μm and 11.2 μm polycyclic aromatic hydrocarbon (PAH) features of a large number of stellar sources, planetary nebulae, reflection nebulae, H II regions, and galaxies, obtained with *Infrared Space Observatory* Short Wavelength Spectrometer. Clear variations are present in the profiles of these features. Most of the sources show a symmetric 3.3 μm feature peaking at $\sim 3.290 \mu\text{m}$, while only very few show an asymmetric 3.3 μm feature peaking at a slightly longer wavelength. The profiles of the 11.2 μm feature are distinctly asymmetric. The majority of the sources has a 11.2 μm feature peaking between 11.20 and 11.24 μm , with a very steep blue rise and a low tail-to-top ratio. A few sources show a 11.2 μm feature with a peak position of $\sim 11.25 \mu\text{m}$, a less steep blue rise, and a high tail-to-top ratio. The sources are classified independently on the basis of the 3.3 and 11.2 μm feature profiles and peak positions. Correlations between these classes and those based on the 6–9 μm features (Peeters et al.) are found. In particular, sources with the most common profiles in the 6–9 μm region also show the most common 3.3 and 11.2 μm feature profiles. However, the uncommon profiles do not correlate with each other. Also, these classifications depend on the type of object. In general, H II regions, nonisolated Herbig AeBe stars and young stellar objects show the same profiles for all 3–12 μm features. Many planetary nebulae and post-asymptotic giant branch stars show uncommon feature profiles. The three galaxies in our sample show the same profiles as the H II regions for all but the 11.2 μm feature, being similar to that of evolved stars. The observed pronounced contrast in the spectral variations for the CH modes (3.3 and 11.2 μm bands) versus the CC modes (6.2, 7.7, and 8.6 μm bands) is striking: the peak wavelengths of the features attributed to CC modes vary by $\sim 15\text{--}80 \text{ cm}^{-1}$, while for the CH modes the variations are $\sim 4\text{--}6.5 \text{ cm}^{-1}$. We summarize existing laboratory data and theoretical calculations of the modes emitting in the 3–12 μm region of PAH molecules and complexes. In contrast to the 6.2 and 7.7 μm components, which are attributed to PAH cations, the 3.3 μm feature appears to originate in neutral and/or negatively charged PAHs. We attribute the variations in peak position and profile of these IR emission features to the composition of the PAH family. The variations in FWHM of the 3.3 μm feature remains an enigma, while those of the 11.2 μm can be explained by anharmonicity and molecular structure. The possible origin of the observed contrast in profile variations between the CH modes and the CC modes is highlighted.

Subject headings: astrochemistry — infrared: ISM — ISM: lines and bands — ISM: molecules — line: identification

1. INTRODUCTION

The infrared spectra of a wide variety of sources are dominated by strong emission features at 3.3, 6.2, 7.7, 8.6, and 11.2 μm (3040, 1615, 1310, 1160, and 890 cm^{-1}), commonly called the unidentified infrared (UIR) emission features (see Gillett et al. 1973; Geballe et al. 1985; Cohen et al. 1986). These features have been detected and studied in a large number of stellar sources—planetary nebulas (PNs), reflection nebulas (RNs), H II regions, young stellar objects (YSOs), and galaxies—indicating that the emitters of these features are surprisingly widespread and extremely stable, and therefore form a very important component of the interstellar medium.

The UIR features are generally attributed to polycyclic aromatic hydrocarbons (PAHs) or PAH-related molecules (Allamandola et al. 1989; Puget & Léger 1989), although the exact molecular identification remains uncertain. Compelling evidence suggests that the features are due to a complex mixture of ionized and neutral, possibly substituted/complexed, PAHs of many different sizes (Schutte et al. 1993; Bakes & Tielens 1994; Boulanger et al. 1998; Bakes et al. 2001; Hony et al. 2001, hereafter HVP01; Verstraete et al. 2001; Pech et al. 2002; Peeters et al. 2002a, hereafter PHV02; Hudgins & Allamandola 2003). Therefore, the overall appearance of the features including the profiles, relative strengths, and peak positions are determined by a large number of PAH-related parameters, which are determined by local conditions and the PAH history.

High-resolution spectroscopy, as obtained with the *Infrared Space Observatory* (ISO; Kessler et al. 1996) and with current ground-based IR instruments, has revealed a richness of widespread variations in the relative strength and profiles of the features from source to source and within sources (Verstraete et al. 1996; Tielens et al. 1999; Maillard et al. 1999; Joblin et al. 2000; HVP01; Vermeij et al. 2002; PVH02; Goto et al. 2003; Madden et al. 2003; Song et al. 2003a, 2003b; Bregman & Temi 2004; Peeters et al. 2004a, 2004b; F. Galliano et al. 2004,

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in preparation). PHV02 have classified the sources in their sample according to the peak positions of the CC modes at 6.2, 7.7, and 8.6 μm independently (Fig. 17 of PHV02 and Figs. 4*b* and 4*c* of this paper). Interestingly, correlations in the features' profiles were found, and three main classes with significantly different feature profiles were identified. Most sources in their sample are in class A (A_{6-9} hereafter). Sources with features peaking at significantly longer wavelengths were referred to as class B (B_{6-9} hereafter). The two sources in their sample showing a 6.3 μm feature exhibit neither a 7.7 μm complex nor a 8.6 μm feature. Instead, both sources show a broad emission feature at 8.22 μm . These sources are referred to as class C (C_{6-9} hereafter). Furthermore, the observed 6–9 μm PAH spectrum was found to depend on the type of object and linked to local physical conditions.

To extend our insight of the PAH family and its relationship with the local physical conditions, we investigated the profiles of the CH modes at 3.3 and 11.2 μm in the *ISO* Short Wavelength Spectrometer (SWS) spectra of a wide variety of sources. In § 2 our sample and the observations are presented; the data reduction, the decomposition of the spectra, and the influence of extinction are discussed. In § 3 the 3.3 and 11.2 μm features are first classified according to their profile and the relations between the classifications of all major features in the 3–12 μm range are investigated, as well as their correlation with object type. Subsequently, variations in features associated with CH and CC modes are compared. In § 4 the spectral characteristics of PAHs in the 3–12 μm range as measured in the laboratory and calculated by quantum chemical theories are summarized. The implications of the spectroscopic, and observational conclusions made in this paper are discussed in § 5. Finally, in § 6 our main results are summarized.

2. THE DATA

2.1. The Sample

For this study, we take the sample of PHV02, which included 57 different objects, all showing UIR bands, and reduce it to those spectra with sufficient signal-to-noise ratio (S/N) in the 3.3 and/or 11.2 μm regions (Table 1). In addition, we excluded HD 100546, which exhibits crystalline silicates (Malfait et al. 1998), and hence the 11.2 μm emission band is a combination of PAH emission and crystalline silicate emission. This reduced sample includes 49 sources from a wide variety of objects, ranging from RNs, (compact) H II regions, YSOs, post-asymptotic giant branch (post-AGB) stars, and PNs to galaxies.

2.2. Observations and Reduction

All spectra in the sample were obtained with the SWS (de Graauw et al. 1996) on board *ISO*, using the AOT01 scanning mode at various speeds or the AOT06 mode, with a resolving power ($\lambda/\Delta\lambda$) of ~ 500 –1500. The data were processed with the SWS Interactive Analysis package IA³ (de Graauw et al. 1996) using calibration files and procedures equivalent with pipeline version 7.0 or later. Further data processing consisted of bad data removal and rebinning with a constant resolution, as described in Peeters et al. (2002b).

In case of high fluxes, the spectra can suffer from memory effects. Here, this only applies to the 11.2 μm feature. At the time the data reduction was done, no memory correction tool was available and the average of the up- and down-scans is taken. The influences of memory effects are investigated as

described in PHV02 and are found not to alter the 11.2 μm feature profile significantly.

2.3. The Continuum

To study the feature profiles in different sources, we subtract a local continuum. This continuum determination is somewhat arbitrary. A local spline continuum or a polynomial of order 2 is fitted to the spectra in a wide wavelength region around the features (Fig. 1). In the 3.3 μm region, the resulting plateau underlying the 3.3 μm PAH band, from about 3.2 to 3.6 μm , was subtracted after the subtraction of the continuum, using a Gaussian with peak position 3.42 μm and a FWHM of 0.242 μm . For the 11.2 μm feature, the plateau and the continuum are fitted together using a single spline. For a review on the plateaus that underly many of the UIR bands, see, e.g., Peeters et al. (2004a). To assess the sensitivity of the resulting profiles to the continuum or plateau choice, two extreme fits were defined and subtracted. In general, the influence of the continuum and/or plateau determination on the profile is very small and hence does not change the classification/profiles of the sources.

2.4. Extinction

The data are not corrected for extinction. The influence of extinction on the 3.3 and 11.2 μm feature profiles is investigated as described in PHV02. The extinction is found to be quite gray over the short wavelength regions and thus has no significant influence on the profiles of the features.

Seven sources show water ice absorption around 3 μm (see Table 1). The 3.3 μm feature is situated in the red wing of the water band, and hence its profile could be influenced. By drawing a local spline continuum (Fig. 1), we compensate to some extent the effect of the ice absorption, and in the following we assume that the ice absorption has no influence on the profile that we derive. This is corroborated by the fact that in our analysis those sources that exhibit clear ice absorption do not show any systematic deviations in the 3.3 μm emission band.

3. THE FEATURE PROFILES

3.1. The 3.3 μm Feature

Nearly all sources in our sample show a pronounced 3.3 μm feature. The similarity of most observed 3.3 μm feature profiles is very striking. The top panel in Figure 2 shows this profile for two quite different object types, the Orion bar H2S1 (H II regions) and NGC 7027 (PN). Only six sources in our sample show deviations in their 3.3 μm features (see Table 1 and Fig. 2, *middle and bottom*).

Based on the profile of the 3.3 μm feature, the sources can be classified into three classes, designated as $A_{3,3}$, $B_{1,3,3}$ and $B_{2,3,3}$ (Tables 1 and 2). Class $A_{3,3}$, containing most sources, has a symmetric profile with a peak position of ~ 3.290 μm and a FWHM of 0.040 μm (Fig. 2, *top*). Classes $B_{1,3,3}$ and $B_{2,3,3}$ have asymmetric profiles with peak positions of ~ 3.293 and ~ 3.297 μm , respectively, both with a FWHM of 0.037 μm . The profile of class $B_{2,3,3}$, represented by HD 44179, is redshifted with respect to class $A_{3,3}$ and has a less steep blue wing (Fig. 2, *bottom*). The profile of class $B_{1,3,3}$, represented by the Orion bar D2, is intermediate; its red wing is similar to the $A_{3,3}$ profile, but the blue wing is clearly shifted to longer wavelengths and the peak position is slightly to the red of the $A_{3,3}$ profile (Fig. 2, *middle*). The up- and down-scans of HR 4049 and IRAS 13416–6243 disagree around the peak,

TABLE 1
THE CLASSIFICATIONS OF THE 3–12 μm PAH FEATURES ASSOCIATED WITH EACH OBJECT

SOURCE	PAH FEATURES					OBJECT TYPE
	3.3 ^a	6.2 ^{a,b}	7.7 ^{a,b}	8.6 ^{a,b}	11.2 ^a	
IRAS 10589–6034 ^{c?}	A	A	A	A	A	Compact H II region
IRAS 12063–6259 ^c	A	A	A	A	A	Compact H II region
IRAS 15384–5348 ^c	A	A	A	A	A	Compact H II region
IRAS 15502–5302 ^c	A	A	A	y	y	Compact H II region
IRAS 18116–1646 ^{c?}	A	A	A	A	A	Compact H II region
IRAS 18032–2032 ^c	A	A	A	A	A	Compact H II region
IRAS 18317–0757 ^c	A	A	A	A	A	Compact H II region
IRAS 18434–0242 ^{c,d}	A	A	A	c	A	Compact H II region
IRAS 19442+2427 ^{c,d}	A	A	A	A	A	Compact H II region
IRAS 21190+5140	A	A	A	A	A	Compact H II region
IRAS 22308+5812	A	A	A	A	A	Compact H II region
IRAS 23030+5958	A	A	A	A	A	Compact H II region
IRAS 23133+6050	A	A	A	A	A	Compact H II region
W3A 02219+6125 ^{c,d}	A	A	A	A	A	Compact H II region
Orion bar D2	B1	A	A	A	A	H II region
Orion bar D5	A	A	A	A	A	H II region
Orion bar D8	A	A	A	A	A	H II region
Orion bar H2S1	A	A	A	A	A	H II region
Orion bar BRGA	A	A	A	A	A	H II region
Orion PK1 ^{c,f}	A	A	A	A	A	H II region
Orion PK2 ^{c,f}	A	A	A	A	A	H II region
G327.3–0.5 ^c	A	A/B1	A	A	A	H II region
IRAS 02575+6017 ^{c,f}	A	A	A	A	A	H II region/YSO
GGD-27 ILL ^{c,f}	A ^g	A	A	A	A	Star-forming region
S106 (IRS4) ^{c,f?}	A	A	A	A	A	YSO
IRAS 03260+3111	A	A	A	A	A	Herbig AeBe star
BD +40°4124	A	A	A	A	A	Herbig AeBe star
CD –42°11721 ^{c?}	A	A	A	A	A	Herbig AeBe star
CD –42°11721(off) ^{c?}	A	A	A	A	A	Herbig AeBe star
HD 97048 ^d	A	A	AB	A	A	Herbig AeBe star
HD 179218	B ^g	B1/2	B	B	y	Isolated Herbig AeBe star
NGC 7023 I	A	A	A	A	A	RN
NGC 2023	A	A	A	A	A	RN
MWC 922 ^d	B1	B1	A	c	y	Emission-line star
CRL 2688	A	C	C	C	y	Post-AGB star
HD 44179 ^d	B2	B2	B	B	B	Post-AGB star
HR 4049 ^h	B	B2	B	B	B	Post-AGB star
IRAS 16279–4757	A	A	A	A	A	Post-AGB star
IRAS 13416–6243 ^h	B	C	C	C	y	Post-AGB star
IRAS 21282+5050	A	A	AB	A	A(B)	PN
BD +30 3639 ^d	A	B1	B	B	A(B)	PN
Hb5	A	A	i	A	B	PN
He 2–113	A	B2	B	B	B	PN
IRAS 07027–7943	A	B2/3	B	B	y	PN
IRAS 17047–5650 ^d	A	B1	B	B	B	PN
NGC 7027 ^d	A	A	B	B	A(B)	PN
Circinus ^c	y	A	A	A	B	Seyfert 2 galaxy
M82 ^{c?,f}	A	A	A	A	A(B)	Starburst galaxy
NGC 253 ^f	A	A	A	A	A(B)	Seyfert galaxy

NOTE.—Question marks in the first column indicate uncertainty; y = feature is present but too weak to classify.

^a Central wavelength is given in μm .

^b From PHV02.

^c Silicate absorption (9.7 μm) present.

^d Sources suffering from memory effects at 11 μm .

^e Source with an unusual 8.6 μm feature (see PHV02; Peeters et al. 2004c).

^f Water ice absorption around 3 μm present.

^g Classification by Van Kerckhoven (2002).

^h Up- and down-scans disagree around 3.3 μm feature, but profile clearly of class B.

ⁱ Strong [Ne VI] present on top of the 7.6 μm feature hampers the classification.

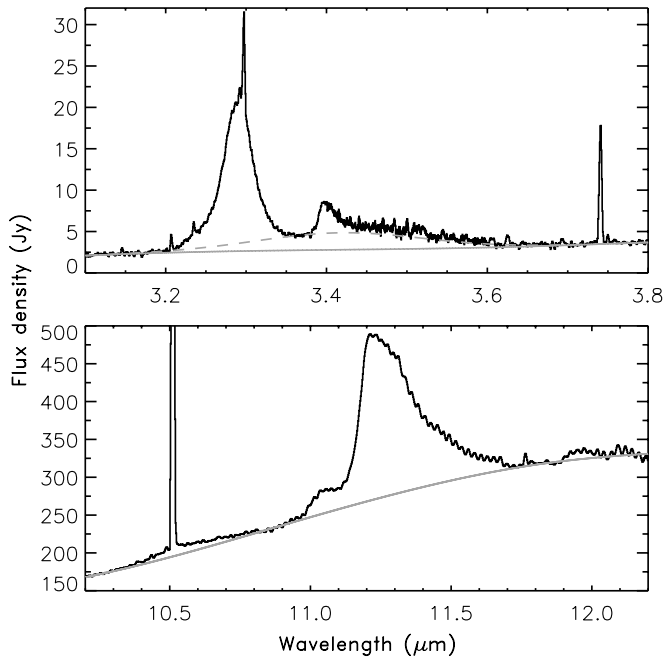


FIG. 1.—Illustrative examples of the continuum underneath the 3.3 and 11.2 μm features exemplified by NGC 7027. The full line represents the local continuum, and the dashed line in the top panel represents the continuum after removing the plateau. The low-amplitude, high-frequency modulations between 11.2 and 12.2 μm are an instrumental artifact due to interference. See § 2.3 for details.

and the spectrum of HD 179218 is noisy. Therefore, we cannot distinguish between classes $B_{1,3,3}$ or $B_{2,3,3}$. However, we can say their profiles do not resemble the class $A_{3,3}$ profile but belong to either $B_{1,3,3}$ or $B_{2,3,3}$, indicated as $B_{3,3}$.

3.2. The 11.2 μm Feature

As a consequence of the selection requirements in § 2.1, the 11.2 μm band profile is not studied for the class C_{6-9} sources and only five sources of class B_{6-9} are included in this sample (Table 1).

All sources in this reduced sample show a pronounced 11.2 μm feature, sometimes preceded by a weak feature at about 11.0 μm . The emission profiles of the 11.2 μm feature are distinctly asymmetric, with a steep blue rise and a red tail. A definite range in peak positions and in the strength of the red wing when normalized so that the profile's peak intensity equals 1 (denoted as tail-to-top ratio) is present in our sample. Perusing the derived profiles, we recognize two main classes, which we will designate by $A_{11,2}$ and $B_{11,2}$ (Table 2, Fig. 3). First, the majority of the 11.2 μm bands peak between 11.20 and 11.24 μm show a very steep blue rise and have a low tail-to-top ratio (Fig. 3, *top*, *black profile*). This group will be referred to as class $A_{11,2}$. Within this class, the peak position and the top of the profiles vary slightly. Class $B_{11,2}$ represents sources with a peak position of ~ 11.25 μm , a less steep blue rise (compared to class $A_{11,2}$) and a high tail-to-top ratio (Fig. 3, *bottom*). A few sources [class $A(B)_{11,2}$] clearly belong to class $A_{11,2}$ in terms of peak position and steepness of the blue wing, but have a tail-to-top ratio similar to class $B_{11,2}$ (Fig. 3, *top*, *gray profile*).

3.3. Feature Profile Correlations

Comparing the classifications of the main PAH bands (Table 1), we find (1) 34 (of 49) sources where all classified

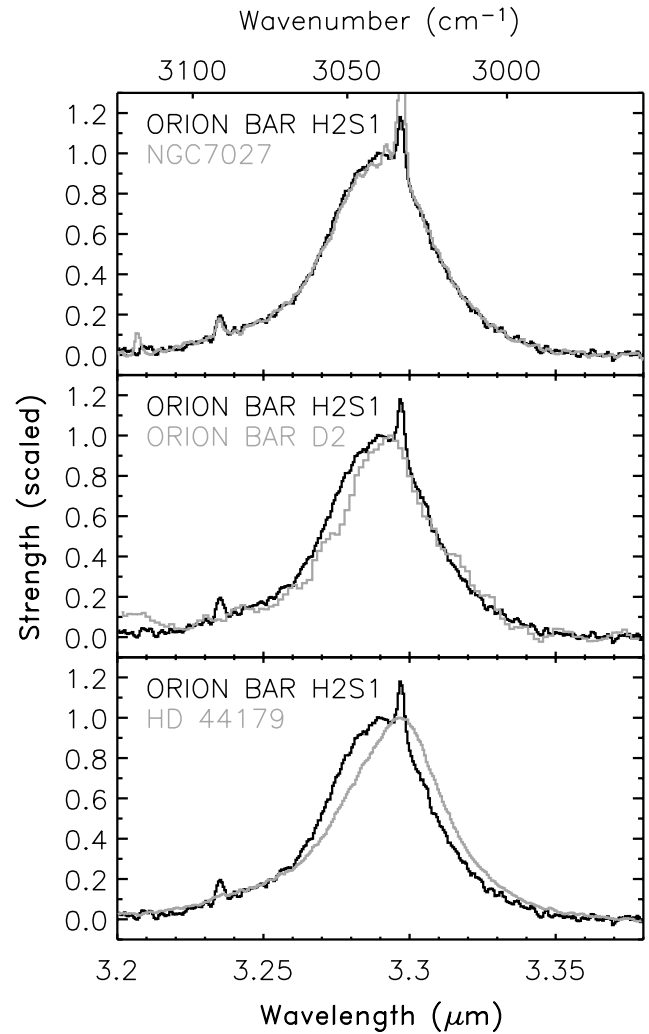


FIG. 2.—Some examples of the 3.3 μm feature. The features are normalized to the peak intensity. Note that a H recombination line (P δ at 3.296 μm) is present in the spectrum of the Orion bar H2S1 and NGC 7027. The top panel shows the shape of the class $A_{3,3}$ features for two quite different sources. This profile is observed in most sources. The middle and bottom panels show the difference between class $A_{3,3}$ and $B_{1,3,3}$, respectively (represented by Orion bar D2) and $B_{2,3,3}$ (represented by HD 44179).

features belong to either class A or class B or class C, of which 31 are class A and three are class B; (2) six sources where all classified features except the 3.3 μm feature belong to either class A or class B or class C; and (3) two sources where all classified features except the 11.2 μm feature belong to either class A or class B or class C. Concluding, we can see that in this large sample of sources, the longer wavelength feature profiles correlate slightly better with each other than with the 3.3 μm feature.

To illustrate the observed variations in the class B profile correlations, a few examples can be given. HD 44179 and HR 4049 have class B profiles, considering all features. BD +30°3639, however, has 6–9 μm features of class B and 3.3 and 11.2 μm features of class A. Other sources, as IRAS 17047–5650 and He 2–113, have 3.3 μm features of class $A_{3,3}$, but class B for the longer wavelength features. Of the two sources of class C_{6-9} , one (IRAS 13416–6243) has a 3.3 μm feature of class $A_{3,3}$ and the other (CRL 2688) of class $B_{3,3}$. In contrast, within the 6–9 μm range, spectral class $B_{7,7}$ also (almost) invariably implies $B_{6,2}$ and $B_{8,6}$ and vice versa.

TABLE 2
OVERVIEW OF THE DEFINED CLASSIFICATIONS OF THE 3.3 AND 11.2 μm PAH FEATURES; SEE TEXT FOR DETAILS

CLASS	PEAK WAVELENGTH		FWHM (μm)	PROFILE
	μm	cm^{-1}		
A _{3.3}	~3.290	~3039	~0.040	Symmetric
B _{13.3}	~3.293	~3037	~0.037	Asymmetric
B _{23.3}	~3.297	~3033	~0.037	Asymmetric
A _{11.2}	~11.20–11.24	~889.6–892.9	~0.17	Asymmetric
A(B) _{11.2}	~11.20–11.24	~889.6–892.9	~0.21	Asymmetric
B _{11.2}	~11.25	~888.9	~0.20	Asymmetric

The 11.2 μm classes follow, in this respect, more the 6–9 μm trend than the 3.3 μm behavior.

Perusing the spectra, it is striking that the variations in the 6–9 μm region are much more pronounced than those of the 3.3 or 11.2 μm features (see Fig. 4; Tables 1 and 3 of PHV02). Specifically, the maximum peak position separation in wavenumber space varies in going from class A to B by about 20, 80, and 15 cm^{-1} for the 6.2, 7.7, and 8.6 μm features, respectively, and only by about 6.5 and 4.0 cm^{-1} for the features at 3.3 and 11.2 μm , respectively.

From an astronomical point of view, correlations between the classes and the different object types are also found. Almost all H II regions, nonisolated Herbig AeBe stars and YSOs in our sample are in class A, considering all features. An interesting exception is the Orion bar position D2, showing a B_{13.3} feature and class A 6–12 μm features. The galaxies in our sample are in class A, considering all features except for the 11.2 μm feature, which is of class A(B)_{11.2} or B_{11.2}. Most PNs have 6–9 μm features of class B_{6–9}, a 3.3 μm feature of class A_{3.3} and a 11.2 μm feature of class B_{11.2} or A(B)_{11.2}. The PN Hb5 has a 11.2 μm feature of class B_{11.2}, while the other bands belong to class A. The post-AGB stars in our sample are spread over the classes. IRAS 16279–4757 is in class A considering all features. IRAS 21282+5050 has a 11.2 μm feature of class B_{11.2}, while the other bands belong to class A. The two extreme metal poor binary system post-AGB stars, HD 44179 and HR 4049, are in class B, considering all features. The same is true for the isolated Herbig AeBe star HD 179218, although its 11.2 μm feature is too weak/noisy to be classified. Thus, in conclusion, the correlation between profile class and object type is much tighter in the 6–9 μm region than in the 3.3 and 11.2 μm regions.

4. THE INFRARED EMISSION FEATURES AND PAHS

The 3.3 and 11.2 μm PAH bands arise from the radiative relaxation of CH stretching and CH out-of-plane bending modes of highly vibrationally excited polycyclic aromatic hydrocarbons. This is in contrast with the nature of the 6.2 and 7.7 μm PAH features, which originate from vibrations mainly involving CC stretching motions. Thus, while one would expect strong global correlations among all of the bands that arise from the same family of molecules, when comparing the behavior of bands involving CH modes with those involving CC modes, subtle but important differences in behavior are expected. These differences reflect not only differences in local conditions, but also structural differences within the emitting family itself, and in the case of the 3.3 μm band emission, from a distinctly different subset of the emitting PAH population.

While discussing the dependence of the precise position of any vibrational mode, one should keep in mind that the vast

majority of the laboratory data available to address each of these bands has been measured in absorption, not emission. In the astrophysical case, where PAH spectra are measured in emission instead of absorption, the extent of molecular excitation also influences peak position and profile. The interstellar spectrum arises from the combined emission of a complex mixture of vibrationally excited PAHs. In emission, each individual line should have an approximately 30 cm^{-1} FWHM, consistent with the natural line width expected from each emitting molecule (Allamandola et al. 1989). This natural line width arises from intramolecular vibrational energy redistribution, not the blending of individual rotational lines in the emitting molecules, which remain rotationally cool throughout the excitation/emission process (e.g., Brenner & Barker 1992; Cook & Saykally 1998). Because of the high internal energy content of the emitting molecules, a 10 to as much as 40 cm^{-1} peak position redshift is intrinsic to the emission process depending on the internal excitation and the emitting feature wavelength (Flickinger et al. 1991; Brenner & Barker 1992;

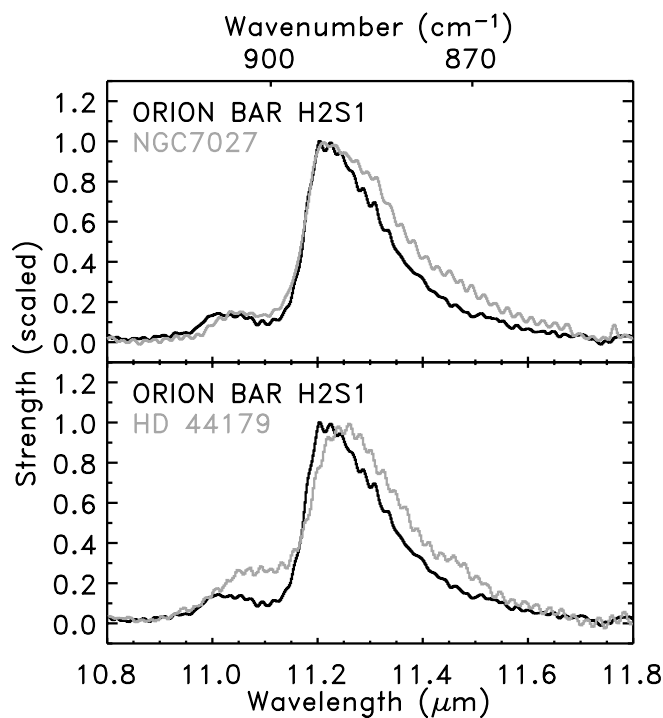


FIG. 3.—Some examples of the 11.2 μm feature. The features are normalized to the peak intensity. The top panel shows the class A_{11.2} feature, represented by the Orion bar H2S1 and the class A(B)_{11.2} feature, represented by NGC 7027. The difference between class A_{11.2} and B_{11.2} (represented by HD 44179) is illustrated in the bottom panel.

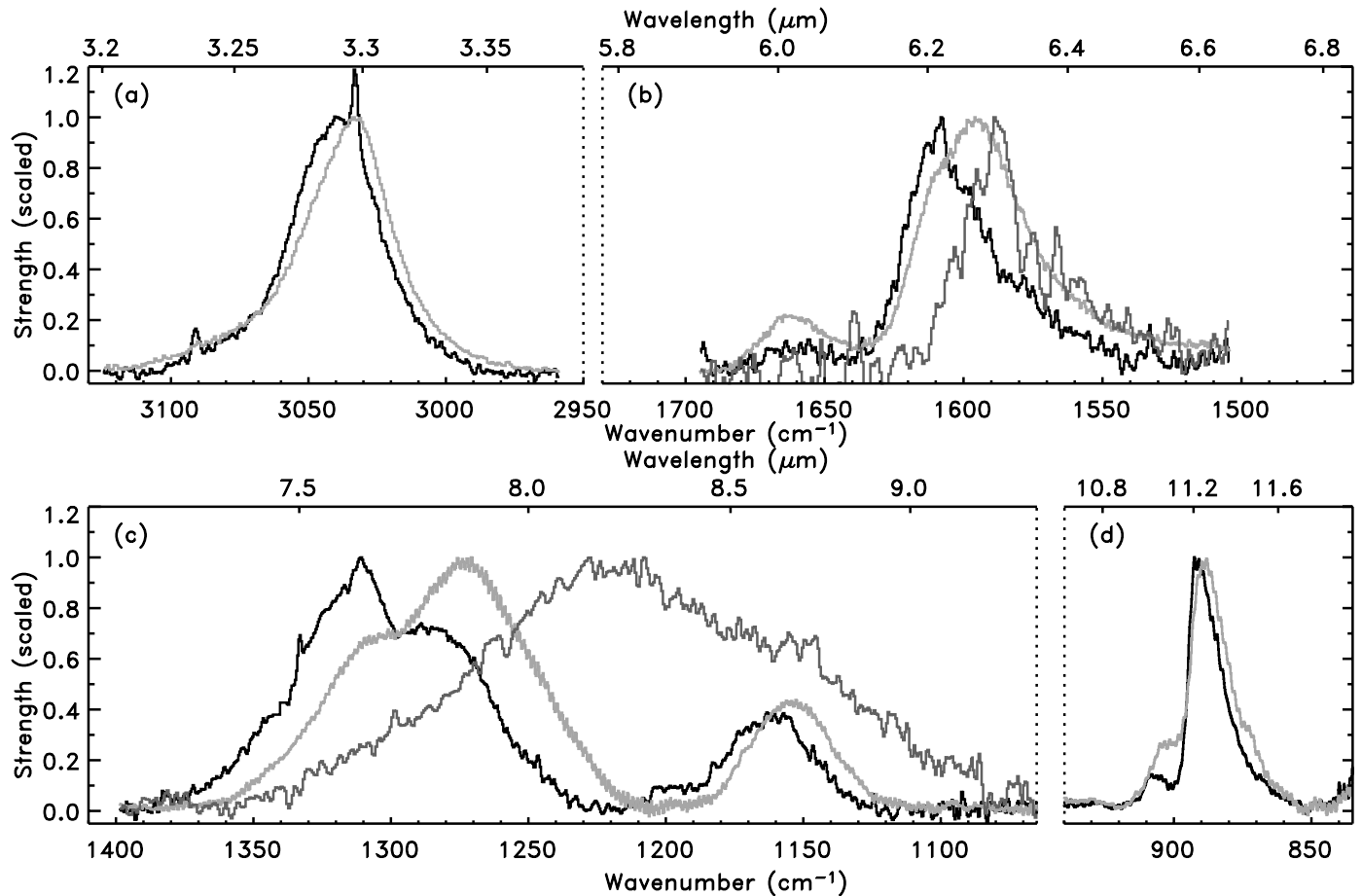


FIG. 4.—Overview of the observed variations in the profiles of the 3–12 μm PAH features. The spectra in each panel are normalized independently to the peak intensity. For all features, class A peaks at the shortest wavelengths (*black line*) and class B peaks at longer wavelength. In the 6–9 μm region, another class (C) is defined, peaking at even longer wavelength and having a deviating spectral appearance (see § 1 and PHV02 for details). Class A is represented by IRAS 23133+6050 in the 6–9 μm region and by the Orion bar H2S1 for the 3.3 and 11.2 μm features. HD 44179 represents class B (classes B_{2,3} and B2 for the 3.3 and 6.2 PAH bands, respectively, and class B for the 7.7, 8.6, and 11.2 μm PAH bands), and IRAS 13416–6243 represents class C. Note that H recombination lines (Pf δ at 3.296 μm and H ϵ (6–5) at 7.458 μm) are present in the spectrum of the Orion bar H2S1 and IRAS 23133+6050.

Colangeli et al. 1992; Joblin et al. 1995; Cook et al. 1998; Wagner et al. 2000).

4.1. The CH Stretching Mode and the 3.3 μm Band

As with all fundamental vibrational frequencies, the precise position of the 3.3 μm band depends on many factors, including charge state, molecular size, structure, molecular heterogeneity, and so on. Each of these factors which influence fundamental band position are discussed below.

Charge state.—The strongest constraint concerns the charge state of the PAHs responsible for the interstellar 3.3 μm emission band. In the Astrochemistry Laboratory at NASA Ames, we now have the experimental infrared spectra of over 100 isolated PAHs in their neutral as well as positively charged states. This spectral database is comprised of PAHs with a number of carbon atoms, N_C , between 10 and 50. These data are supplemented with theoretical spectra that extend the size to $N_C = 96$. The number of PAHs for which anion spectra are available is considerably smaller, currently 27. Analysis of these spectra shows that the CH stretching frequencies are sensitive to PAH charge as well as size (D. M. Hudgins et al. 2004, in preparation). This is illustrated in Figure 5, which compares the interstellar emission feature from NGC 7027 (class A_{3,3}) to the charge-specific CH stretch regions for the entire Ames sample. Figure 5 shows

that the CH stretch in PAH anions spans the range from 3085 to 3025 cm^{-1} (3.24–3.30 μm), neutral PAHs from 3094 to 3042 cm^{-1} (3.23–3.29 μm), and PAH cations from 3112 to 3073 cm^{-1} (3.21–3.25 μm). The symbol within each range indicates the average position for the PAHs in the Ames database. These are 3048, 3081, 3063 and 3093 cm^{-1} (3.28, 3.24, 3.26 and 3.23 μm) for the anionic forms, large neutrals, small neutrals, and cationic forms, respectively.

PAH cations seem least likely to contribute to the 3.3 μm feature. First, Figure 5 shows that they require the largest redshift possible (40 cm^{-1}) to overlap only the blue wing of the feature. Second, PAH cations have inherent weak emission in the 3.3 μm region, i.e., the CH band of cations is suppressed by a factor 50–100 compared to that of neutral PAHs (Langhoff 1996). Lastly, the 3.3 and 11.2 bands are well correlated and behave differently than the 6.2 and 7.7 bands that are dominated by PAH cations (Hony et al. 2001). In contrast, Figure 5 shows that the CH stretches for neutral PAHs, when redshifted as a result of the high internal energy content, fall directly under the most intense portion of the interstellar 3.3 μm band envelope, while the PAH anions require a small redshift to fall underneath the 3.3 μm band envelope. The breadth and small variance of the interstellar 3.3 μm emission band likely arises because the CH stretching position (in absorption) for the vast majority of isolated, neutral and anionic PAHs falls between

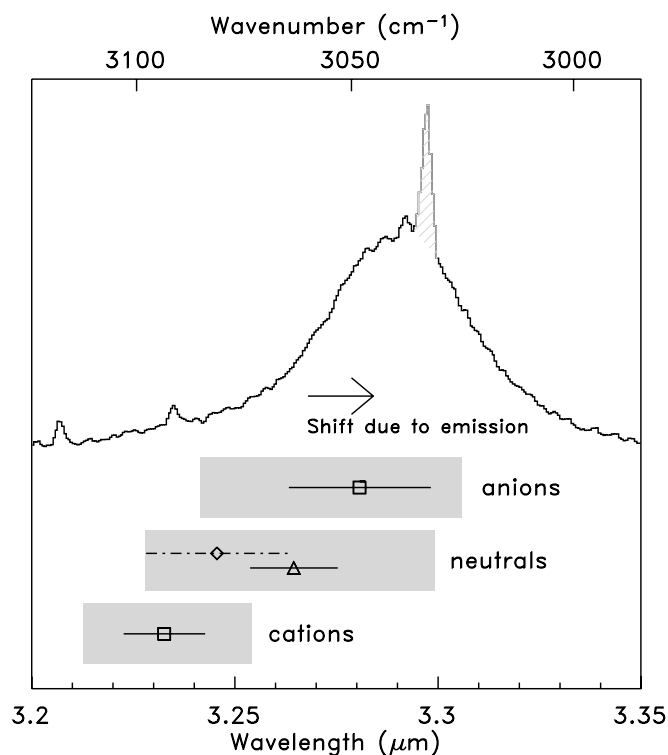


FIG. 5.—Comparison of the interstellar emission feature from NGC 7027 (class A_{3.3}) with the charge-specific CH stretch regions for the entire Ames sample. The striped, gray feature is Pf β . The shaded areas represent, for each charge state, the full measured range in positions for the CH stretch mode. The symbols crossed by horizontal lines represent the average position and its standard deviation. The diamond and triangle crossed by horizontal lines are specific for large ($N_C \geq 40$) and small ($N_C < 40$) neutrals, respectively. The arrow represents a 15 cm^{-1} peak position redshift intrinsic to the emission process (see text for details).

3085 to 3025 cm^{-1} . With the 10 – 40 cm^{-1} redshift inherent in emission, this range becomes 3075 – 2985 cm^{-1} (3.25 – $3.35 \mu\text{m}$), a range which straddles the wavelength region of this feature listed in Table 2.

Size.—PAH size is particularly important in the case of the interstellar $3.3 \mu\text{m}$ band since this feature originates from the smallest members of the emitting PAH population (Allamandola et al. 1989; Schutte et al. 1993). PAHs containing between roughly 25 – 70 carbon atoms contribute most of the emission in this feature, with the species N_C between ~ 25 and 50 dominant. This is in contrast to the situation for the longer wavelength emission features, which arise from increasingly larger and overlapping members of the interstellar PAH population.

Interestingly, within the size range of the current Ames sample, there appears to be a bimodal distribution of the frequencies for the neutral species (see Fig. 5). For the neutral PAHs, the CH stretch for the smallest members of the sample is centered near 3060 cm^{-1} ($3.27 \mu\text{m}$), while the larger members have frequencies centered at 3090 cm^{-1} ($3.24 \mu\text{m}$). The origin of the bimodal distribution is not clear from the current sample.

Edge structure.—PAH edge structure is also important. If the periphery of an individual PAH is not regular, the local environment that each CH bond samples can perturb the individual stretching frequencies. Thus, for a large molecule with long straight edges, there is little interaction between the stretching motions of CH groups on adjacent rings and the

FWHM of the dominant peak in the region is 18 cm^{-1} ($0.019 \mu\text{m}$) at $3.3 \mu\text{m}$. On the other hand, if a similarly large molecule that has jagged edges with many “notches” and “bay regions,” CH groups on opposing sides of these gaps interact, perturbing their stretching motions. As a result, the FWHM of the CH stretching feature in such a species is substantially broadened and may be as wide as 70 cm^{-1} ($0.074 \mu\text{m}$) or more (D. M. Hudgins et al. 2004, in preparation).

Heteroatom substitution.—Mattioda et al. (2003) have investigated the spectra of 11 N-substituted PAHs, spanning the N_C range from 9 to 21, and found that nitrogen substitution does not influence the CH stretching band frequency. N substitution does influence other PAH vibrational transitions (PHV02; Mattioda et al. 2003; D. M. Hudgins et al. 2004, in preparation)

4.2. The CH Out-of-Plane Bending Modes

The aromatic CH out-of-plane bending features in the 11 – $15 \mu\text{m}$ spectral region are a good diagnostic for the classification of the aromatic ring edge structures. Indeed, the positions and intensities of the bands in this spectral region reflect the type and number of adjacent CH groups on the peripheral rings of the PAH structure (Bellamy 1958; Allamandola et al. 1985; Cohen et al. 1985; Léger et al. 1989; Roche et al. 1989; Witteborn et al. 1989; Allamandola et al. 1999; Hudgins & Allamandola 1999, HVP01). Our current understanding is illustrated in Figure 6, which schematically compares the average interstellar emission spectrum with the wavelength regions associated with different CH adjacency classes for neutral and ionized, isolated PAHs. Inspection of this figure

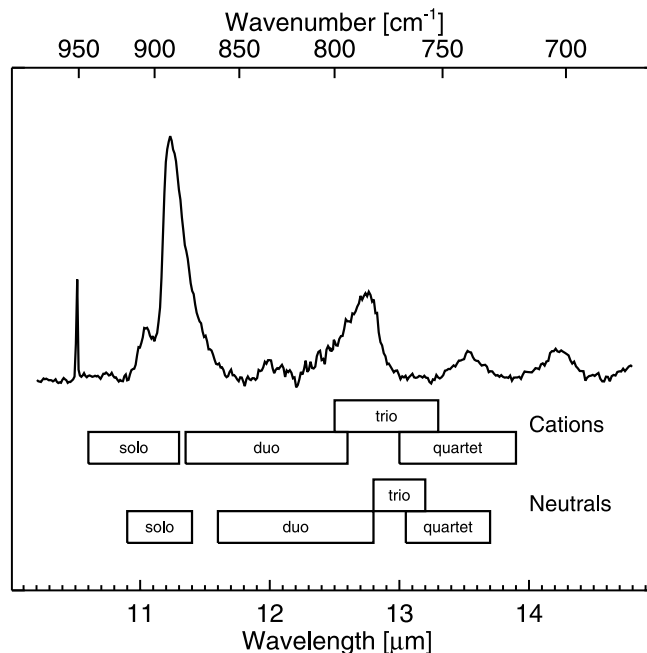


FIG. 6.—Comparison of the average interstellar spectrum (top) with the ranges characteristic of the CH out-of-plane bending modes of neutral and cationic PAHs (bottom). Details of the average interstellar spectrum are given in HVP01. The boxes indicate the wavelength regions associated with the out-of-plane bending vibrations of the various adjacency classes of the peripheral CH groups as determined from matrix isolation spectroscopy. Aromatic rings carrying CH groups that have no neighboring CH groups are termed “non-adjacent” or “solo” CH groups. Likewise, two adjacent CH groups are called “duet” CHs, three adjacent CH groups “trio” CHs and four adjacent CH groups “quartet” CHs. Figure taken from HVP01.

shows that, while the ranges for matrix-isolated neutral PAHs do not differ substantially from those reported in the literature, ionization produces some notable changes in region boundaries (Hudgins & Allamandola 1999). Taking into account the roughly 0.1 μm redshift in the peak wavelength of 11.2 μm PAH band in emission relative to their position in absorption, the following conclusions are drawn (Hudgins & Allamandola 1999; HVP01): (1) The weak interstellar emission band sometimes observed on the short-wavelength side of the dominant 11.2 μm band and peaking near 11.0 μm falls in the region attributable to the solo-CH modes of PAH cations. (2) The peak of the 11.2 μm interstellar band falls squarely within the region for the solo-CH modes of *neutral* PAHs and at the extreme long-wavelength end of the range for solo cationic modes. (3) The variable red wing of the interstellar 11.2 μm band likely carries a contribution from the duet-CH modes of PAH cations, particularly those with condensed structures whose features tend to fall near the lower extreme of the characteristic region shown in Figure 6.

4.3. The CC Stretching Modes

The CC stretching modes occur in the 6–9 μm range. Within this range, the bands at the longer wavelengths have a mixed CC stretch and CH in-plane-bending vibration character. All CC modes are very weak relative to the CH modes in neutrals but become the dominant bands in ions (Szczepanski & Vala 1993; Langhoff 1996; Hudgins et al. 1994; Kim et al. 2001). Hence, we will focus here on the CC modes in cations. The peak position of the pure CC stretch is influenced by molecular size, structure, and molecular heterogeneity (see, e.g., PHV02). Extensive laboratory and theoretical quantum mechanical studies have shown that a peak as blue as 6.2 μm does not occur for pure-C PAH cations in the (interstellar) size range 20–100 C atoms. Substitution of N deep in the fused ring system induces strong IR activity of the CC modes at the highest frequencies within this range, and bands around 6.2 μm do become intense (PHV02; D. M. Hudgins et al. 2004, in preparation). This is the result of the delocalization of the charge across the entire PAH as a result of incorporating N in the ring. While there is presently no corroborating experimental or theoretical evidence, a similar behavior may occur for large PAHs with uneven and irregular edge structures, PAH clusters, and/or PAH complexes with metals. For the longer wavelength modes, both pure C PAH cations as well as N-substituted PAH cations can show prominent bands near the interstellar 7.6 μm position. The origin of the 7.8 μm band is more enigmatic, however, and no good match with laboratory species exists.

5. DISCUSSION

From the wealth of medium resolution IR spectra, it is clear that the major PAH features show clear variations in profile and peak position. Large variations are found in profiles of the features in the 6–9 μm region, while smaller ones are found in those of the 3.3 and 11.2 μm features. Furthermore, the 3.3 μm feature is very invariant under highly variant physical conditions. The different profile classes are linked with each other and with object type; this correlation, however, is more apparent in the 6–9 μm region than in the 3.3 and 11.2 μm regions.

5.1. The Observed Trends

Tokunaga et al. (1991) recognized two types of 3.3 μm features based on spectra taken with a 2''7 aperture. The pro-

file, peak position, and width of their type 1 agree well with class A_{3,3} defined here. Their type 2 has a peak wavelength similar to that of our classes B_{3,3} but is much narrower, with a FWHM of $\sim 0.020 \mu\text{m}$. Most sources of their sample belong to type 1, while type 2 is associated with only three sources (HD 44179, Elias 1, WL 16). The 3.3 μm PAH band in the ISO/SWS spectra (obtained through an aperture of 14'' \times 20'') of their type 1 sources belongs to class A_{3,3} while those of HD 44179 and Elias 1 belongs to class B_{3,3} (see Table 1 and Van Kerckhoven [2002] for Elias 1). Kerr et al. (1999), and Song et al. (2003a) showed, in a detailed spatial study of HD 44179, that the 3.3 μm profile varies, going from type 2 to type 1 with distance from the central star. Hence, the observed profile will depend on the sampled area on the sky. In particular, the 3.3 μm profile of HD 44179 in our sample is observed through an aperture 14'' \times 20'' and thus has an averaged profile of both type 1 and type 2 profiles, thereby explaining its large FWHM compared to the profile found by Tokunaga et al. (1991), Kerr et al. (1999) and Song et al. (2003a). In addition, the type 1/class A_{3,3} profile is not unique to certain object types. Whether the type 2 profile only occurs in specific sources is currently unknown. However, few sources are found that show a type 2 or class B_{3,3} profile. All except one are post-AGB stars (three), isolated HAeBe stars (three), and an emission-line star (Molster et al. [2002] classifies MWC 922 as disklike based on the basis of the spectral energy distribution and its mm continuum flux). Apparently, a circumstellar environment where dust formation is ongoing or a protostellar/protoplanetary environment is an essential but not sufficient condition to have type 2 or class B_{3,3} profiles. Indeed, some post-AGB stars or HAeBe stars show a type 1/class A_{3,3} profile. The fact that the 3.3 μm feature varies considerably less than the other main features and mostly exhibits the A_{3,3}/type 1 profile indicates that the conditions that give rise to the B_{3,3} and/or type 2 profiles are more easily lost than those that give rise to the B (and C) types for the other bands. Song et al. (2003b) and Miyata et al. (2004) also reported a varying 8.6 and 11.2 μm profile in HD 44179, with a peak position going to shorter wavelengths with distance from the star and hence going from class B to class A. Note also that the 3.3 μm feature also shifts blueward with distance from the star, indicating that these three PAH feature are correlated in this source. As with the 3.3 μm feature, the observed profile will depend on the sampled area on the sky. In addition, class A profile is not unique to certain object types.

Tokunaga et al. (1991) noted that the central stars of the few sources showing a type 2 profile, all had relatively low effective temperatures ($\sim 10^4$ K), suggesting a correlation between stellar effective temperature and PAH feature profiles. However, both profile types are found in HD 44179 (Kerr et al. 1999; Song et al. 2003a). At the same time, some objects in our sample have a relatively low stellar temperature, e.g., IRAS 03260+3111 (Jaschek et al. 1964), but still show an A_{3,3} class profile. This demonstrates that the effective temperature cannot be the sole cause of the variations in the 3.3 μm feature profile. Van Kerckhoven (2002) and PHV02 have reached similar conclusions for the longer wavelength features.

Since variations are seen within objects, it may be that the local radiation field, G_0 , influences the profiles of the feature. However, no link between G_0 and the different classes of the 3.3 and 11.2 μm features is found. Indeed, for both features individually, the range in G_0 of the sample sources classified as class A straddles that of class B and class A(B)_{11.2} sources. Analogously, PHV02 found no relation between G_0 and

profiles of the 6–9 μm PAH emission features. Hence, we can conclude that G_0 alone cannot be the factor determining the observed variation in the main PAH emission features from source to source (observed through a large aperture).

5.2. The Origins of the Variations in the 3.3 and 11.2 μm CH Emission Features

In the framework of the PAH hypothesis, the profiles and thus also the variations in the profiles reflect the present molecular ensemble and the excitation conditions. In the following, we examine whether we can isolate any effect which is clearly reflected in the observed variations of the profiles.

5.2.1. Anharmonicity

The profiles are very characteristic for anharmonic broadening in highly vibrationally excited molecules (Barker et al. 1987; Kim et al. 2001; Verstraete et al. 2001; Pech et al. 2002). Because of cross-coupling between different modes, the peak position of an IR band will shift toward lower frequencies. Integration over the energy cascade leads then to a red shaded profile. The red wing can be further enhanced by emission from vibrationally excited levels (e.g., vibrational hot bands, $v = 2 \rightarrow 1$), which are generally shifted toward the red as a result of the anharmonicity of the vibrational potential of the mode. These anharmonicities can be quite different for different modes. Anharmonicity parameters for a few small PAHs have been estimated from laboratory studies on the temperature dependence of the peak wavelength and width of the absorption bands (Joblin et al. 1995) and from the difference between low-temperature absorption measurements and infrared multiphoton dissociation studies of trapped PAH ions (Oomens et al. 2003).

The effects of anharmonicity have been studied theoretically (Verstraete et al. 2001; Pech et al. 2002). These studies show that the general characteristics of the interstellar 3.3, 6.2, and 11.2 μm bands are well understood in the framework of highly excited PAHs. However, these models cannot explain the large shift observed in the 6.2 μm band between classes A, B, and C. Such a shift can only reflect a chemical/structural modification of the emitting PAHs, for example, due to the incorporation of N (PHV02; D. M. Hudgins et al. 2004, in preparation).

The calculated 3.3 μm profiles of Pech et al. (2002) show that an increase in excitation—either through a decreased size or an increase in internal energy—shifts the peak position toward the red and simultaneously increases the FWHM (note that the calculated 3.3 μm profile has a larger asymmetry factor than observed in space). Our class $B_{3,3}$ profiles and the type 2 profiles peak at longer wavelengths but have smaller FWHM than class $A_{3,3}$ /type 1 profiles. Thus, the variation in profiles between classes $A_{3,3}$ /type 1 and either $B_{3,3}$ or type 2 cannot be due to anharmonicity.

Perusing the Pech et al. (2002) 11.2 μm model profiles, we notice that an increase in excitation—either through a decreased size or an increase in internal energy—does increase the strength of the red wing but cannot shift the peak position or change the blue rise. In their study of the 11.2 μm band, Pech et al. (2002) focused on IRAS 21282+5050, a class $A(B)_{11,2}$ source, and indeed the $A(B)_{11,2}$ profile can be interpreted as originating from the same PAH population that gives rise to the $A_{11,2}$ profile excited to higher internal energy. In contrast, the difference between $A_{11,2}$ and $B_{11,2}$, with the band

as a whole displaced, are not readily explained within the anharmonicity model.

5.2.2. Composition

The 3.3 and 11.2 μm CH bands are likely due to the presence of a family of PAH molecules, and hence the variation in peak position and profile of the CH modes likely represents a small but significant difference in the population of emitting species. The variation in the peak position of the 3.3 and 11.2 μm band, measured in the laboratory or quantum-chemically calculated, is much larger (0.02 μm , 20 cm^{-1} and 0.3 μm , 20 cm^{-1} , respectively) than observed in space (~ 0.007 μm , 6 cm^{-1} and ~ 0.05 μm , 4 cm^{-1} , respectively). Although the convolution of the size distribution with the excitation conditions can lead to a different dispersion for the peak positions (compared to absorption bands measured in the lab), it is likely that only a *small subset* of the PAHs studied by these means contribute to the observed emission from space.

5.2.3. Molecular Edge Structure

The 3.3 μm feature.—If we assume that the variation in the FWHM of the 3.3 μm feature between class $A_{3,3}$ /type 1 and either $B_{3,3}$ or type 2 is mainly influenced by the molecular edge structure, this observed variation in FWHM (and hence in class) implies that the PAH family present in class $B_{3,3}$ and type 2 sources are dominated by more compact PAH molecules than the PAH family present in class $A_{3,3}$ /type 1 sources (§ 4.1 and Table 2). Since the molecular structure of the PAH molecules can be deduced on the basis of the relative strengths of the CH out-of-plane bending modes (HVP01), and assuming that the PAHs responsible for the 3.3 μm band have similar molecular structure as those for the 11.2 μm band (although these features arise from the emission of a different subset of PAHs), we can test the dependence of the FWHM of the 3.3 μm feature on molecular structure. Let us consider the following sources: HD 44179 belonging to $B_{3,3}$ and $B_{11,2}$; NGC 7027, member of $A_{3,3}$ and $B_{11,2}$; and the IRAS 18317, representing classes $A_{3,3}$ and $A_{11,2}$. HVP01 found that NGC 7027 is dominated by large, compact PAHs with straight edges and IRAS 18317 by smaller or irregular PAHs. Although these two sources represent the extremes in the observed intensity ratio in their sample, and hence the extremes in molecular structure, they both belong to $A_{3,3}$. In addition, the relative strengths of the CH out-of-plane bending modes of HD 44179 are comparable to NGC 7027 and hence their molecular structure is similar (HVP01). However, their 3.3 μm features belong to different classes. Therefore, we can conclude that the molecular edge structure is likely not the main cause of the variation in the FWHM of the 3.3 μm feature.

The 11.2 μm feature.—The little variance of the CH out-of-plane bending modes may be an inherent property of interstellar PAHs. Within the subset of PAHs studied in the laboratory or quantum chemically, the peak position of the solo modes in neutral and cation PAHs falls very closely to the observed interstellar peak position of 11.2 μm , with the neutral PAHs dominating the emission (see Fig. 8 in HVP01). Now, possibly, any molecular structure variations, which lead to shifts in the 6–9 μm CC bands, may be inherently accompanied by modifications of the CH peripheral structure that turn CH solos into duos or trios.

Duet CHs of compact PAH cations may contribute to the red wing of the 11.2 μm band (Hudgins & Allamandola 1999, § 4.1; HVP01). Hence, it is possible that the observed

variation in the 11.2 μm band reflects variation in the contribution of these cationic duet modes. This would imply that $B_{11.2}$ and $A(B)_{11.2}$ objects have more compact PAH cations with duets. Analysis of the pattern of emission features in the CH out-of-plane bending modes region (i.e., 10–15 μm ; HVP01) reveals that evolved stars are dominated by large compact structures while H II regions are dominated by irregular structures and hence is consistent with this conclusion.

5.2.4. Carbon Isotope Effects

On the basis of laboratory studies, shifts in the peak positions of the interstellar IR emission features have been attributed to ^{13}C isotope substitution (Wada et al. 2003). We note that while this might be important in some extreme circumstellar environments (such as the Red Rectangle), this cannot be a general explanation for the observed shifts. First, not all the features shift in the same way at all times. Second, we see pronounced variations in material that is interstellar in origin and where isotope substitution should be minimal. The observed variations within a single H II region or reflection nebula present, of course, is a case in point (C. Joblin et al. 2004, in preparation; Bregman & Temi 2004).

To summarize, the variation in peak position is likely due to the presence of a mixture of PAH molecules whose composition changes. Both anharmonicity and molecular structure can explain variation in the FWHM of the 11.2 μm feature. In contrast, the variation in the FWHM of the 3.3 μm feature between class $A_{3.3}$ /type 1 and type 2 remains an enigma, although it might be the result of a change in composition of the PAH family. The larger FWHM of our class $B_{3.3}$ likely originates from the average (through a large aperture) of type 2 and type 1/class $A_{3.3}$ profiles with a dominating type 2 contribution.

5.3. The Different Behavior of the CH Modes versus the CC Modes

Perusing the spectra, it is striking that the variations in the 6–9 μm region are much more pronounced than those of the 3.3 or 11.2 μm features (see Fig 4; Table 2). This implies that variations in the features assigned to CC modes are significantly larger than those assigned to CH modes. Here we discuss what could be the root of this dichotomy.

5.3.1. Charge State

Experimentally and quantum chemically, it is well established that the 6–9 μm CC modes are dominated by cations while the 3.3 μm CH stretching mode can be due to neutrals and anions (Szczepanski & Vala 1993, Fig. 5 and § 4.1; Langhoff 1996; Hudgins & Allamandola 1999; Kim et al. 2001; Mattioda et al. 2004, in preparation). There is some observational support that this dichotomy between the CC and CH stretching modes also extends to the 11.2 μm CH out-of-plane bending mode (HVP01). Specifically, the strength of the astronomical 11.2 μm mode correlates well with the 3.3 μm CH stretching band and not with the 6.2 μm CC stretching band. The difference in spectral variations between CH and CC modes may merely reflect that they probe different parts of the interstellar PAH population. For example, the chemical modification process might be mainly operative on cations and not on neutrals. In that respect, because ionization and neutralization of interstellar PAHs occur on a very rapid time-

scale, this would imply a similar fast chemical exchange. Such a fast process would also allow for spatial variations within a source. It should be noted that it is unlikely that this rapid behavior is connected with the incorporation of N in PAH rings. However, functionalized PAHs, PAH clusters, or PAH-metal complexes may be formed and broken on a similar timescale (Van Kerckhoven 2002).

5.3.2. Heteroatom Substitution and Complexes

Substitution of nitrogen in the ring can explain variation in the 6.2 μm band position. However, it does not influence the peak position of the CH modes, nor does it influence the peak positions of the 7.7 μm complex in a systematic way. However, as mentioned above, it is unlikely that N incorporation deep within a PAH can occur on such a rapid timescale as to explain the spatial variations within a source. Complexes of PAHs with each other or with metal atoms might also cause the spectral variation in the 6–9 μm range. However, the influence of complexation on the CH modes still has to be investigated in the laboratory.

5.3.3. Size

The very small variation of the 3.3 μm feature with respect to the CC modes and to the CH out-of-plane bending modes is likely also related to distinct *subsets* of emitting PAH molecules. Indeed, the 3.3 μm feature originates from the smallest members of the emitting PAH population, while the PAH features in the 6–12 μm region originate from somewhat larger PAHs in the population (Schutte et al. 1993). This would indicate that photochemistry is not the driving force behind the observed spectral variation since smaller PAHs will be much more susceptible to unimolecular dissociation following photon absorption and hence larger variations would be expected for the 3.3 μm feature than for higher wavelength features, in contradiction to what is observed. In addition, this might suggest that while the larger PAHs are modified, the smaller ones are destroyed. This would leave only the most stable small PAHs to emit the 3.3 μm feature, while a mixture of modified larger PAHs survive to give rise to the higher wavelength features.

5.3.4. Molecular Edge Structure

As already mentioned, any molecular structure variations that lead to shifts in the 6–9 μm CC bands may be inherently accompanied by modification of the CH peripheral structure which turn CH solos into duos or trios.

To summarize, several parameters can cause the difference in the amount of variation between the CH and CC modes. However, further laboratory studies are required to determine which of these parameters(s) controls the spectral characteristics of PAHs in space. Based on the rough correlation of most of the band profiles, one can stipulate that the appearance of the UIR emission features is determined by a few physical parameters alone. However, there are several sources that do not follow these correlations. This implies that the behavior cannot be understood in terms of a single parameter.

6. CONCLUSIONS

Based on *ISO/SWS* observations, we have studied the profiles of the 3.3 and 11.2 μm PAH features of a wide variety of sources and found that clear variations are present. Both features shift in peak position and show different profiles from

source to source. Comparing the 3.3, 6–9, and 11.2 μm classifications, we recognize a correlation between them. In general, an A classification in 6–9 μm region also implies an A_{3.3} and A_{11.2}. However, this seems not true for the B and C classifications; B_{3.3} and B_{11.2} do not necessarily correlate with each other or with B_{6–9} or C_{6–9}. In addition, these variations depend on the type of object considered. Apparently, a circumstellar environment where dust formation is ongoing or a protostellar/protoplanetary environment is an essential but not sufficient condition to have peculiar 3.3 or 11.2 μm profiles. Noteworthy is the fact that the three galaxies in our sample show the same profile as the H II regions, except for the 11.2 μm profile, which resembles that of some evolved stars.

Spatial variations of the profiles within a source indicate that specific profiles are not unique to certain object types and that the observed profile depends on the aperture. Nevertheless, the average (or predominant) profile (as observed with *ISO/SWS*, which has an aperture of $14'' \times 20''$) present within a source does depend on object type.

However, the most striking aspect of the features in the 3–12 μm regions is the pronounced contrast in the profile variations between the CH modes and the CC modes. Specifically, the peak position in wavenumber space varies by about 20, 80, and 15 cm^{-1} for the 6.2, 7.7, and 8.6 μm features, respectively, and only by about 6.5 and 4.0 cm^{-1} for the features at 3.3 and 11.2 μm , respectively.

We summarize existing laboratory data and theoretical calculations of the modes emitting in the 3–12 μm region of PAH molecules. In contrast to the 6–9 μm region, which arises from PAH cations, the 3.3 μm band appears to originate mainly in neutral and/or negatively charged PAHs.

We attribute the variations in peak position of the 3.3 and 11.2 μm feature to the presence of a mixture of PAH molecules, whose composition changes. The variations in FWHM of the 3.3 μm feature remains an enigma although they might result from the change in composition of the PAH family—while those of the 11.2 μm can be explained by anharmonicity and molecular structure. The possible origin of the observed contrast in profile variations between the CH modes and the CC modes is highlighted.

The overall good agreement that can be achieved between the IS spectra and the spectral characteristics of free PAHs strongly supports the PAH model and allows these spectra to become a probe of local conditions. However, it is clear that some features of the major band profiles and the minor bands cannot be accounted for by “classical” PAHs. Other PAH related species, such as PAH clusters, heteroatom PAHs, and PAH-metal complexes, should also be considered.

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